# An experimental overview of effective field theory exploration at the LHC

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### A few other caveats...

- This talk will cover several analyses
- I will be focusing more on the interpretation of the analyses as opposed to specific selection criteria used in an analysis
- ✓ I have focussed on recent results that use the full LHC Run II dataset
- Mostly focused on CMS results



# The Standard Model of Particle Physics But is it complete?



#### **Open questions:**

- Origin of neutrino mass?
- What is Dark matter?
- Not a complete theory that includes gravitational interactions

Focus of today's talk: bosons ¥, W, Z

# **Standard Model Measurements**



### Text book example of effective field theories



### **Higher Dimensional Operators**



### **Effective field theories**

### Complementary to BSM searches



Instead of bump hunting, looking for new physics in the tails of distributions

# The EFT framework is useful an example from the leptonic sector



- Tension in the leptonic sector with the SM at 3.6  $\sigma$  observed in  $b o s l^+ l^-$
- EFT fits performed that factor 3.1  $\sigma$  deviation from SM in measurement of  $R_k$
- Combination of these measurements leads to a 4.6  $\sigma$  significance of the hypothesis of a purely left-handed lepton flavor universality-violating contact interaction

$$R_K = \frac{\mathcal{B}(B^+ \to K^+ \mu^+ \mu^-)}{\mathcal{B}(B^+ \to K^+ e^+ e^-)}$$

 $R_{K^{(*)}}$ 

0.0

 $\Delta C_{0}^{\mu}$ 

0.5

1.0

# Search for new physics with multibosons at the multipurpose detectors



# **Muon detection in CMS**





### **Processes of interest**

#### **Diboson production at the LHC**





### W<sup>±</sup> γ production

### W+ W- /WZ production

In all cases, decay at least one gauge boson leptonically

# **Exploration of dimension-6 operators**

$$\mathcal{L} = \mathcal{L}_{\mathcal{SM}} + \sum_{i} rac{c_i}{\lambda^2} \mathcal{O}_i + \sum_{j} rac{f_j}{\lambda^4} \mathcal{O}_j + \cdots$$

Use Standard Model Effective Field Theory or SMEFT framework to characterize impact of dim-6 operators

$$\sigma(C_{3W}) = \sigma_{SM} + C_{3W}\sigma_{\text{interference}} + C_{3W}^2\sigma_{BSM}$$
  
Linear Quadratic

# **Slew of dimension-6 operators**

<b>1</b> : <i>X</i> <sup>3</sup>		2 : <i>H</i> <sup>6</sup>			3 : H <sup>4</sup> D <sup>2</sup>	<b>5</b> : $\psi^2 H^3$ + h.c.			
$Q_G$	$f^{ABC}G^{A u}_\mu G^{B ho}_ u G^{C\mu}_ ho$	Q <sub>H</sub> (	$(H^{\dagger}H)^{3}$	$Q_{H\Box}$	$(H^{\dagger}H)\Box(H^{\dagger}H)$	Q <sub>eH</sub>	$(H^{\dagger}H)(\bar{l}_{p}e_{r}H)$		
$Q_{\widetilde{G}}$	$f^{ABC}\widetilde{G}^{A u}_{\mu}G^{B ho}_{ u}G^{C\mu}_{ ho}$			Q <sub>HD</sub>	$\left(H^{\dagger} \mathcal{D}_{\mu} H\right)^{*} \left(H^{\dagger} \mathcal{D}_{\mu} H\right)$	$Q_{uH}$	$(H^{\dagger}H)(\bar{q}_{\rho}u_{r}\widetilde{H})$		
$Q_W$	$\epsilon^{I\!J\!K} W^{I u}_\mu W^{J ho}_ u W^{K\mu}_ ho$					$Q_{dH}$	$(H^{\dagger}H)(\bar{q}_{p}d_{r}H)$		
$Q_{\widetilde{W}}$	$\epsilon^{I\!J\!K}\widetilde{W}_{\!\mu}^{I\! u}W_{\! u}^{J ho}W_{\! ho}^{K\mu}$								
<b>4</b> : <i>X</i> <sup>2</sup> <i>H</i> <sup>2</sup>		(	<b>δ</b> : ψ <sup>2</sup> ΧΗ -	+ h.c.		<b>7</b> : ψ <sup>2</sup> <i>H</i> <sup>2</sup> <i>D</i>			
$Q_{HG}$	$H^{\dagger}HG^{A}_{\mu u}G^{A\mu u}$	Q <sub>eW</sub>	$(\bar{l}_{p}\sigma^{\mu u}e$	$(e_r)  au^I H W^I_{\mu u}$	$Q_{HI}^{(1)}$	(H <sup>†</sup> i <sup>ć</sup> D	$\partial_{\mu}H)(\bar{l}_{\rho}\gamma^{\mu}l_{r})$		
$Q_{H\widetilde{G}}$	$H^{\dagger}H\widetilde{G}^{\mathcal{A}}_{\mu u}G^{\mathcal{A}\mu u}$	$Q_{eB}$	$(ar{l}_{ m p}\sigma^{\mu u}$	$(e_r)HB_{\mu u}$	$Q_{HI}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D})$	$(\bar{l}_{ m p}  au^{\prime} \gamma^{\mu} l_{ m r})$		
$Q_{HW}$	$H^{\dagger}HW^{I}_{\mu u}W^{I\mu u}$	$Q_{uG}$	$(ar{q}_{ ho}\sigma^{\mu u}$ 7	$(\mathcal{F}^{A}u_{r})\widetilde{H}G^{A}_{\mu\nu}$	v Q <sub>He</sub>	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}_{p}\gamma^{\mu}e_{r})$			
$Q_{H\widetilde{W}}$	$H^{\dagger}H\widetilde{W}^{I}_{\mu u}W^{I\mu u}$	$Q_{uW}$	$(ar{q}_{ m p}\sigma^{\mu u})$	$(J_r)  au' \widetilde{H} W'_{\mu\nu}$	$, \qquad Q_{Hq}^{(1)}$	$  (H^{\dagger}i\overleftrightarrow{D})$	$_{\mu}H)(ar{q}_{ ho}\gamma^{\mu}q_{r})$		
$Q_{HB}$	$H^{\dagger}HB_{\mu u}B^{\mu u}$	$Q_{uB}$	$(ar{q}_{ ho}\sigma^{\mu u}$	$(u_r)\widetilde{H}B_{\mu u}$	$Q_{Hq}^{(3)}$	$  (H^{\dagger}i\overleftrightarrow{D}^{I}_{\mu})$	$(H)(\bar{q}_{p} au^{\prime}\gamma^{\mu}q_{r})$		
$Q_{H\widetilde{B}}$	$H^{\dagger}H\widetilde{B}_{\mu u}B^{\mu u}$	$Q_{dG}$	$(ar{q}_{ ho}\sigma^{\mu u}$ 7	$(^{A}d_{r})HG^{A}_{\mu\nu}$	ν Q <sub>Hu</sub>	$  (H^{\dagger}i\overleftrightarrow{D})$	$_{\mu}H)(\bar{u}_{p}\gamma^{\mu}u_{r})$		
$Q_{HWB}$	$H^{\dagger} au^{\prime}HW^{\prime}_{\mu u}B^{\mu u}$	$Q_{dW}$	$(ar{q}_{ m p}\sigma^{\mu u}c)$	$(d_r) \tau^I H W^I_{\mu\nu}$	, Q <sub>Hd</sub>	$  (H^{\dagger}i\overleftrightarrow{D})$	$_{\mu}H)(ar{d}_{ ho}\gamma^{\mu}d_{r})$		
$Q_{H\widetilde{W}B}$	$H^{\dagger} au^{\prime}H\widetilde{W}^{\prime}_{\mu u}B^{\mu u}$	$Q_{dB}$	$(ar{q}_{ m p}\sigma^{\mu u}$	$(d_r)HB_{\mu u}$	Q <sub>Hud</sub> + h.c.	$i(\widetilde{H}^{\dagger}D_{\mu})$	$_{\iota}H)(\bar{u}_{p}\gamma^{\mu}d_{r})$		

### **Description of analyses - dim6 exploration**







EFT exploration performed with combination of four final states:
 WW and WZ in leptonic final states
 Final state consisting of four charged leptons
 Leptonically decaying Z-boson in vector-boson-fusion topology

### W±γ CMS: <u>https://arxiv.org/abs/2111.13948</u>



# Effective field theory constraints from the W±γ process



Focus is on

$$\mathcal{O}_{3W} = \epsilon^{abc} W^{a\nu}_{\mu} W^{b\rho}_{\nu} W^{a\mu}_{\rho}$$

 $\sigma(C_{3W}) = \sigma_{SM} + C_{3W}\sigma_{\text{interference}} + C_{3W}^2\sigma_{BSM}$ 

- The interference with the SM cannot be experimentally detected with "inclusive" observables (decay angles) for *ff* → W<sub>T</sub>V<sub>T</sub>
- Different helicity configurations for SM and BSM components
- Leads to suppression of interference effects
- Can probe only pure-BSM contributions
- "Diboson interference resurrection" performed (<u>arXiv:1708.07823</u>, <u>arXiv:1901.04821</u>)
  - Requires definition of specific coordinate system

### Special coordinate system



### Constraints on Wilson Coefficients from the W±γ process



## Measurement of the WW + 1 jet







Specific focus on production W+W- with an additional jet

Access to trilinear gauge coupling which could acquire contributions from dim-6 operators

Just like in the Wγ analysis, interference with SM suppressed due to helicity configurations in SM vs. BSM

 $\mathcal{N}_{W^+}$ 

q

✓ Production mode also includes loop-induced gluon–gluon fusion process (gg → WW)

> W+W-CMS: Phys. Rev. D 102 092001 (2020) ATLAS: Eur. Phys. J. C 79 (2019) 884 JHEP 06 (2021) 003

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# Measurement of the WW + 1 jet





#### https://arxiv.org/pdf/1707.08060.pdf

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no

# Measurement of the WZ process





Electroweak process: sensitive to the PDFs of u and d quarks; relatively unaffected by the gluon

High WZ cross section makes it the dominant process that can be studied in the trilepton final state

Ratio of W+Z/W-Z cross section is one of the most precisely measurable quantities

Constitutes first measurement of longitudinally polarized W-bosons θ<sup>w</sup>: angular distance between the momenta of the W boson and the charged lepton from its primary decay

#### CMS: https://arxiv.org/abs/2110.11231

# Measurement of the WZ process



Unfolded distributions of sensitive variables well modeledSimulated signal samples normalized to NNLO cross sections

# Measurement of the WZ process



Parameter	95% CI, Exp. (TeV <sup>-2</sup> )	95% CI, Obs. (TeV <sup>-2</sup> )	Best fit, Obs. $(\text{TeV}^{-2})$	
$c_{\rm w}/\Lambda^2$	[-2.05, 1.27]	[-2.52, 0.33]	-1.34	CP conserving
$c_{\rm www}/\Lambda^2$	[-1.27, 1.33]	[-1.04, 1.19]	0.15	
$c_{\rm b}/\Lambda^2$	[-86.0, 125.0]	[-42.7, 113.0]	43.6	CP non-conserving
$\tilde{c}_{ m www}/\Lambda^2$	[-0.76, 0.65]	[-0.62, 0.53]	-0.03	
$\tilde{c}_{\rm w}/\Lambda^2$	[-46.1, 46.1]	[-45.9, 45.9]	0.0	

Limits computed taking interference with SM ( $\Lambda^{-2}$ ) + pure BSM ( $\Lambda^{-4}$ ) into account

**M** Improvement over limits from W+W<sup>-</sup> analysis (taking luminosity scaling into account) for  $C_w$  and  $C_{www}$ 

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Invariant mass of the secondary lepton pair (the

pair that is less compatible with the Z boson mass)

**ATLAS** Preliminary

 $\sqrt{s} = 13 \text{ TeV}, 139 \text{ fb}^{-1}$ 

4l,  $m_{_{4l}} \approx m_{_{Z}}$  region

do/dm<sub>z2</sub> [fb/GeV] Stat. Syst. Uncertainty Stat. Syst. Uncertainty Off-shell region explored in addition to on-shell  $m_{\rm T}^{WZ} = \sqrt{\left(\sum p_{\rm T}^{\ell} + E_{\rm T}^{\rm miss}\right)^2 - \left(\sum \vec{p}_{\rm T}^{\ell} + \vec{E}_{\rm T}^{\rm miss}\right)^2}$ 10-100 600 700 500 900 1000 n 200 300 400 800 10 15 20 25 30 35 40

m<sup>WZ</sup><sub>T</sub> [GeV]

**EFT** exploration performed with combination of four final states:

WW and WZ in leptonic final states

Final state consisting of four charged leptons

Prediction

Theory Uncertainty

Measurement

Leptonically decaying Z-boson in vector-boson-fusion topology

Effect of 33 SMEFT operators explored

**ATLAS** Preliminary

 $\sqrt{s} = 13 \text{ TeV}, 36 \text{ fb}^{-1}$ 

da/dm<sup>WZ</sup> [fb/GeV] 0

WZ

**Combined effective field theory interpretation of** differential cross-sections measurements of WW, WZ, 4I, and Z-plus-two-jets production using ATLAS data

45

m<sub>72</sub> [GeV]

50

ATLAS: http://cdsweb.cern.ch/record/2776648

Prediction

Theory Uncertainty

Measurement

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Combined effective field theory interpretation of differential cross-sections measurements of WW, WZ, 4I, and Z-plus-two-jets production using ATLAS data

Form of the gauge boson operators under exploration (leptonic operators also explored the analysis, focusing only on a subset)

Wilson	coefficient and operator	Final state affected at leading order									
	_	$e^{\pm}\nu\mu^{\mp}\nu$	$\ell^+\ell^-\ell^\pm \nu$	4 <i>ℓ</i>	$\ell^+\ell^- jj$						
$c_G$	$f^{abc}G^{a u}_{\mu}G^{b ho}_{\nu}G^{c\mu}_{ ho}$				$\checkmark$						
$c_W$	$\epsilon^{IJK}W^{I u}_{\mu}W^{J ho}_{ u}W^{K\mu}_{ ho}$	$\checkmark$	$\checkmark$		$\checkmark$						
$C_{HD}$	$\left(H^{\dagger}D_{\mu}H ight)^{*}\left(H^{\dagger}D_{\mu}H ight)$		$\checkmark$	$\checkmark$	$\checkmark$						
$C_{HWB}$	$H^\dagger  au^I H  W^I_{\mu u} B^{\mu u}$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
$\overline{c_{Hl}^{(1)}}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{l}\gamma^{\mu}l)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
$c_{Hl}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{l}\tau^{I}\gamma^{\mu}l)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
C <sub>He</sub>	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{e}\gamma^{\mu}e)$		$\checkmark$	$\checkmark$	$\checkmark$						
$c_{Ha}^{(1)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{q}\gamma^{\mu}q)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
$c_{Hq}^{(3)}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}^{I}H)(\bar{q}\tau^{I}\gamma^{\mu}q)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
$C_{Hu}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{u}\gamma^{\mu}u)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						
$C_{Hd}$	$(H^{\dagger}i\overleftrightarrow{D}_{\mu}H)(\bar{d}\gamma^{\mu}d)$	$\checkmark$	$\checkmark$	$\checkmark$	$\checkmark$						



#### **Combined effective field theory interpretation of** differential cross-sections measurements of WW, WZ, 4I, and

z-pi	us-two-jets production using ATLAS dat		ERIMENT
Liı	inear Effect of Wilson Coefficient		
Ex	xperimental Uncertainty		
	ATLAS Preliminary	Linear Effect of Wils	son Coefficient
-	√s = 13 TeV, 36-139 fb <sup>-1</sup>	Lin+Quad Effect of	Wilson Coefficient
	Relative Effect of Wilson Coefficient, for $\Lambda = 1$ TeV	Experimental Unce	rtainty
	0.4 0.3 0.2 0.2		
	27 >300 0 >600 5 50 5 9	955 90 1100-π m [GoV]	π Λά
	$p_{T}$ [Gev] $p_{T} \rightarrow ev_{\mu\nu}$ $p_{T} \rightarrow III_{\nu}$ $p_{T} \rightarrow II_{\nu}$ $p_{T} \rightarrow II_{\nu}$	$p \rightarrow 4l$ pp	→ IIjj
	Impact of the Wilson coefficie	ents on the	
and the second secon	differential cross sections, rela	ative to the SM	

cross section

ΛC

Value of the Wilson coefficient obtained from 20 sensitivity of the SMEFT model including only linear terms

m<sup>WZ</sup> [GeV]

 $pp \rightarrow III_V$ 

>6005 505

≃m<sub>z</sub> off-shell

955

 $pp \rightarrow 4l$ 

m<sub>4</sub>>2m<sub>z</sub> 90 1100-π

m<sub>z2</sub> [GeV]

ATLAS Preliminary √s = 13 TeV, 36-139 fb<sup>-1</sup>

<sup>0.4</sup><sub>0.3</sub> c<sub>w</sub>=0.29

c<sub>HD</sub>=0.44

с<sub>нwв</sub>=0.21

c<sub>HI</sub><sup>(1)</sup>=0.4

с<sup>(3)</sup>=0.4

с<sub>не</sub>=0.54

с<sup>(1)</sup>=0.99

с<sup>(3)</sup>=0.11

с<sub>ни</sub>=2.0

c<sub>Hd</sub>=5.1

c<sub>II</sub><sup>(1)</sup>=0.6

 $pp \rightarrow ev\mu v$ 

>300 (  $p_{\tau}^{\text{lead. lep}}$  [GeV]

Relative Effect of Wilson Coefficient, for  $\Lambda = 1 \text{ TeV}$ 

Δφ..

pp → lljj

Combined effective field theory interpretation of differential cross-sections measurements of WW, WZ, 4I, and Z-plus-two-jets production using ATLAS data



**ATLAS** Preliminary  $\sqrt{s} = 13 \text{ TeV}$ , 36-139 fb<sup>-1</sup>

c' <sup>[0]</sup> 4q	0.07	0.14	0.17	0.08	0.18	0.10	-0.06	0.10	-0.13	-0.47	-0.13	-0.17	-0.11	-0.23	1.00		tion 1
C' <sup>[3]</sup> 2q2l	-0.13	-0.85	-0.51	0.06	-0.12	-0.27	-0.45	-0.37	0.35	0.04	0.45	0.71	0.60	1.00	-0.23		orrela
c' <sup>[2]</sup> 2q2l	-0.13	-0.68	-0.19	0.05	-0.46	0.25	-0.68	-0.28	0.49	-0.03	0.47	0.38	1.00	0.60	-0.11		ů Ne
c' <sup>[1]</sup> 2q2l	-0.14	-0.86	-0.41	0.10	0.09	-0.68	-0.50	-0.70	0.11	0.01	0.58	1.00	0.38	0.71	-0.17		0.0
c' <sup>[0]</sup> 2q2l	-0.13	-0.58	-0.20	0.04	-0.20	-0.22	-0.54	-0.61	0.07	0.11	1.00	0.58	0.47	0.45	-0.13		0.4
c' <sup>[7]</sup> Vff	-0.03	0.12	-0.09	-0.12	-0.15	-0.03	0.14	-0.06	-0.02	1.00	0.11	0.01	-0.03	0.04	-0.47		02
c' <sup>[6]</sup> Vff	0.01	-0.33	-0.09	0.03	-0.21	0.17	-0.26	-0.03	1.00	-0.02	0.07	0.11	0.49	0.35	-0.13		0.2
c' <sup>[5]</sup> <sub>Vff</sub>	0.12	0.55	0.23	-0.08	-0.02	0.46	0.40	1.00	-0.03	-0.06	-0.61	-0.70	-0.28	-0.37	0.10	_	0
c' <sup>[4]</sup> Vff	0.13	0.65	0.14	-0.07	0.21	0.05	1.00	0.40	-0.26	0.14	-0.54	-0.50	-0.68	-0.45	-0.06		_02
c' <sup>[3]</sup> Vff	0.04	0.36	0.24	-0.04	-0.37	1.00	0.05	0.46	0.17	-0.03	-0.22	-0.68	0.25	-0.27	0.10		-0.2
c' <sup>[2]</sup> <sub>Vff</sub>	0.08	0.07	0.01	0.05	1.00	-0.37	0.21	-0.02	-0.21	-0.15	-0.20	0.09	-0.46	-0.12	0.18		-0.4
c' <sup>[1]</sup> Vff	-0.00	-0.14	-0.05	1.00	0.05	-0.04	-0.07	-0.08	0.03	-0.12	0.04	0.10	0.05	0.06	0.08		_0.6
c' <sup>[0]</sup>	0.08	0.40	1.00	-0.05	0.01	0.24	0.14	0.23	-0.09	-0.09	-0.20	-0.41	-0.19	-0.51	0.17		0.0
$c_{Hq}^{(3)}$	0.15	1.00	0.40	-0.14	0.07	0.36	0.65	0.55	-0.33	0.12	-0.58	-0.86	-0.68	-0.85	0.14		-0.8
$c_W$	1.00	0.15	0.08	-0.00	0.08	0.04	0.13	0.12	0.01	-0.03	-0.13	-0.14	-0.13	-0.13	0.07		_1
	$c_W$	$c_{Hq}^{\left( 3 ight) }$	$c_{\rm Vff}^{\rm [0]}$	$c_{\rm Vff}^{\rm [1]}$	$c_{\rm Vff}^{\rm [2]}$	$c_{\rm Vff}^{\rm [3]}$	$c_{Vff}^{\rm [4]}$	c' <sup>[5]</sup> Vff	c' <sup>[6]</sup>	$c_{Vff}^{[7]}$	c' <sup>[0]</sup> 2q2l	c' <sup>[1]</sup> 2q2l	c' <sup>[2]</sup> 2q2l	c' <sup>[3]</sup> 2q2l	c' <sup>[0]</sup> 4q		1
Correlation between the 15 parameters of interest																	

- Focus of this analysis is the simultaneous measurement of all Wilson coefficients
- Construct a modified basis (linear combination of Warsaw basis vectors <u>https://</u> <u>cds.cern.ch/record/2694284</u>, <u>https://cds.cern.ch/record/</u> <u>2743067</u>) to reduce flat directions and identify sensitive directions
- Hessian matrix constructed in the space of Wilson coefficients and eigenvectors are identified
- This analysis sets the stage for global combination efforts

### **Processes of interest**

Vector-boson scattering production at the LHC
Triboson production\*





#### \*more on that later

# **Exploration of dimension-8 operators**

$$\mathcal{L} = \mathcal{L}_{\mathcal{SM}} + \sum_{i} rac{c_i}{\lambda^2} \mathcal{O}_i + \sum_{j} rac{f_j}{\lambda^4} \mathcal{O}_j + \cdots$$

**<u>CP</u>** conserving quartic operators are of the form:

$$\mathcal{L}_{\mathrm{S},0} = \left[ (D_{\mu} \Phi)^{\dagger} D_{\nu} \Phi \right] \times \left[ (D_{\mu} \Phi)^{\dagger} D_{\nu} \Phi \right] \qquad \mathcal{L}_{\mathrm{M},0} = Tr \left[ \hat{W}_{\mu\nu} \hat{W}^{\mu\nu} \right] \times \left[ (D_{\beta} \Phi)^{\dagger} D^{\beta} \Phi \right]$$
$$\mathcal{L}_{\mathrm{T},0} = Tr \left[ W_{\mu\nu} W^{\mu\nu} \right] \times Tr \left[ W_{\alpha\beta} W^{\alpha\beta} \right]$$

# Vector boson scattering processes



Final state consists of same-signed leptons and 2 high p<sub>T</sub> jets

#### Rapidity gap between jets (no color flow)

Provides unique access to quartic couplings

Large invariant mass of the two jets

#### Use of Zeppenfeld variable

$$z_l^* = |\eta_l - \frac{1}{2}(\eta_{j1} + \eta_{j2})| / |\Delta \eta_{jj}|$$



# Exploration of WW and WZ processes with two additional jets



CMS: Phys. Lett. B 809 (2020) 135710

# Wy + 2 jets



Exploration of full set of "mixed" operators performed

For the parameters f<sub>M2-5</sub>, f<sub>M6-7</sub> most stringent limits

First observation of the W $\gamma$  + 2 jets process with observed (expected) significance of 5.3 (4.8)  $\sigma$ 

Exploration of dim-8 operators possible due to presence of SM quartic coupling

Invariant mass of the Wγ system is sensitive to presence of dim-8 operators



### Factoring in EFT validity

- EFT validity taken into account by restricting EFT integration at the unitarity limit
- ✓ Unitarity limit is at ~1.5 TeV

#### Philosophy 1. Disregard unitarity limits (CMS mainstream)

- technically simplest,
- fair to quantify the relative precision of different measurements and the degree of agreement/disagreement with the SM,
- obtained numbers do not have direct EFT interpretation and unitarity violation usually occurs well within the measured range.

#### **Philosophy 2. Unitarization techniques:**

#### Amplitude saturation, e.g., K-matrix (ATLAS mainstream), Form factor approach (e.g. VBFNLO)

- describe the maximum possible signal related to a given operator,
- no unique prescription,
- part of a model,
- obtained numbers not easy to interpret within the EFT.

#### Truncation, validity, uncertainties

Ilaria Brivio, Sally Dawson, Jorge de Blas, Gauthier Durieux, Pierre Savard (editors), Roberto Contino, Céline Degrande, Adam Falkowski, Florian Goertz, Christophe Grojean, Fabio Maltoni, Ken Mimasu, Giuliano Panico, Francesco Riva, William Shepherd, Eleni Vryonidou, Andrea Wulzer, Cen Zhang https://arxiv.org/abs/2201.04974v1

#### Limits worse by a factor of 5 in comparison with limits without unitarity constraints but more physical

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From Michal Szleper's talk

## Triboson processes



# **Cross sections X branching fractions**



### WWW production in the Standard Model



**Require:** 

Two leptons of <u>identical charge</u> arising out of the decay of two W-bosons
#### Search for WWW in dilepton final state

- Search for WWW in:
  - 2 same-signed leptons + 2 jets
    - further categorized based on  $|M_{jj} M_W| \le 15 \text{ GeV}$
  - 2 same-signed leptons + 1 jet
- Major backgrounds are WZ and nonprompt contribution, some prompt (W<sup>±</sup>W<sup>±</sup>jj / ttW)
- BDTs trained against nonprompt and prompt backgrounds



#### WWW in dilepton final state

WWW  $\rightarrow$  2 lepton + 2 jet event



CMS experiment at the LHC, CERN CMS

Data recorded: 2016-Jul-02 14:25:40.606976 GMT

Run 276242, Event No. 96020969 LS 52

#### Search for WWZ in four final state



# Require four leptons targeting decay products arising out of the W and Z-bosons

#### Very clean channel!

- Requiring the presence of four leptons leads to the presence of almost no nonprompt leptons
- Further requiring that there be two Zboson candidates, as expected in a ZZ process, can lead to a control region that is 98% pure in ZZ
- This control region enables the study of the ZZ process, which is the major background in this channel
- Need to identify variables that can lead to the discrimination of ZZ from WWZ process (target signal)



# Some interesting variables to separate signal from backgrounds



- Can break down signal region into various components:
  - Require the presence of one Z-boson
    - Further split the signal region based on the presence of e-µ pairs
    - ZZ contributes through  $\tau\tau$  decays (of each Z) leading to e- $\mu$  pairs but mass of the e- $\mu$  pair will be lower than the Z-mass due to presence of neutrinos

#### Search for WZZ in five lepton final state



# Require five leptons targeting decay products arising out of the W and Z-bosons

## WZZ in five lepton final state

- In 5 lepton channel:
  - Require 2 Z boson candidates and associate remaining lepton with a W
  - Separate by flavor of the W candidate lepton and require M<sub>T</sub> > 50 GeV if electron



## WZZ in five lepton final state



#### ZZZ in six lepton final state



## Require six leptons targeting decay products arising out of the Z-bosons

#### Search for ZZZ in six lepton final state

- In 6 lepton channel:
  - Require ∑ pT of all leptons
     > 250 GeV, powerful
     against backgrounds which
     contribute at percent level



#### **Combination of all final states**



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#### **Observation of Heavy Tribosons with Run II data**

Major milestone in Standard Model physics!



#### PHYSICAL REVIEW LETTERS

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Observation of the Production of Three Massive Gauge Bosons at  $\sqrt{s}=13~{
m TeV}$ 

A. M. Sirunyan *et al.* (CMS Collaboration) Phys. Rev. Lett. **125**, 151802 – Published 5 October 2020

Physics See synopsis: Hat Trick Observation for Bosons

Auxiliary material: http://cms-results.web.cern.ch/cms-results/public-results/publications/SMP-19-014/

#### **Observation of Heavy Tribosons with Run II data**

Culmination of persistent endeavor by both <u>ATLAS</u> and <u>CMS</u> Collaborations



Recent <u>result</u> from ATLAS on observation of the WWW process

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#### Effective field theory exploration in WWW process

$$\mathcal{L} = \mathcal{L}_{\mathrm{SM}} + \sum_{i \frac{c_i}{\lambda^2}} \mathcal{O}_i + \sum_{j \frac{f_j}{\lambda^4}} \mathcal{O}_j$$





# $\mathcal{L}_{\mathrm{T},0} = \mathrm{Tr} \left[ W_{\mu\nu} W^{\mu\nu} \right] \times \mathrm{Tr} \left[ W_{\alpha\beta} W^{\alpha\beta} \right]$

#### Effective field theories in VVV process

Currently pursuing an Effective Field Theory interpretation with VVV process

Including as many final states as possible

✓ Presence of higher order operators expected to lead to high p<sub>T</sub> decay products: 2 resolved jets → 1 fatJet

#### Explore dim-6 and dim-8 operators

Process	Final states					
	Fully leptonic	Hadronic decay of one gauge boson	Leptonic decay of one gauge boson	Fully hadronic		
WWW	$3 \ell + \text{missing energy}$	$2 \ell + \text{jets} (W^{\pm} \to \ell^{\pm} \nu, W^{\pm} \to \ell^{\pm} \nu, W \to q\bar{p})$	$1 \ \ell + \text{jets} \ (W^{\pm} \to \ell^{\pm} \nu, W^{\pm} \to q\bar{p}, W \to q\bar{p})$	6 jets		
WWZ	$4 \ell + \text{missing energy}$	$2 \ \ell + \text{jets} \ (W \to \ell \nu, W \to \ell \nu, Z \to q\bar{q})$	$1 \ \ell + \text{jets} \ (W \to \ell \nu, W \to q\bar{p}, Z \to q\bar{q})$	6 jets		
		$3 \ell + \text{jets} (W \to q\bar{p}, W \to \ell\nu, Z \to \ell^+ \ell^-)$	$2 \ell + \text{jets} (W \to q\bar{p}, W \to q\bar{p}, Z \to \ell^+ \ell^-)$			
WZZ	$5 \ell + \text{missing energy}$	$3 \ \ell + \text{jets} \ (W \to \ell \nu, Z \to \ell^- \ell^+, Z \to q\bar{q})$	$1 \ \ell + \text{jets} \ (W \to \ell \nu, Z \to q\bar{q}, Z \to q\bar{q})$	6 jets		
		$4 \ \ell + \text{jets} \ (W \to q\bar{p}, Z \to \ell^- \ell^+, Z \to \ell^- \ell^+)$	$2 \ \ell + \text{jets} \ (W \to q\bar{p}, Z \to \ell^- \ell^+, Z \to q\bar{q})$			
ZZZ	$6\ell$	$4 \ \ell + \text{jets} \ (\mathbf{Z} \to \ell^- \ell^+, \ \mathbf{Z} \to \ell^- \ell^+, \ \mathbf{Z} \to q\bar{q})$	$2 \ \ell + \text{jets} \ (\mathbf{Z} \to q\bar{q},  \mathbf{Z} \to q\bar{q},  \mathbf{Z} \to \ell^- \ell^+)$	6 jets		

Final states explored in observation paper

 $Z \rightarrow \nu \overline{\nu}$  not included for brevity

#### Salient features associated with VVV

- Study of the Higgs sector possible
  - Best way to include Higgs mediated mode?
  - Can we quantify the interference between the Higgs-mediated modes and other modes of production?
- Critical questions such as (dim6)<sup>2</sup> vs. dim8 contributions possible to address
  - Pertinent in the context of global EFT fits
- Explore sensitive variables: proxy for  $\hat{s}$  and angular variables ( $\Delta \phi$ )
  - Absence of "golden" variables as in the case of vector boson scattering topologies



## The LHC upgrade schedule



#### **VVV significance at HL-LHC**

Process	Final State	Significance expected $\sigma$	
WWW	same-signed dilepton and trilepton	Can be studied in the realm of precision physics	
WWZ	four leptons	Can be studied in the realm of precision physics	
WZZ	five leptons	> 5.0	
ZZZ	six or more leptons	> 5.0	

#### This is a simple projection assuming uncertainties remain the same

#### **Power of new detectors** High Granularity Calorimeter

- Unprecedented spatial granularity
- 3D visualization of showers
- Timing capabilities
- Use low shower time resolution (~ few ps) to distinguish between spatially overlapping showers
- Use *ps* timing to mitigate pileup







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#### Path forward...

- Presented an exploration of new physics in diboson, vector boson scattering (VBS), triboson topologies using the effective field theory framework presented
- Elucidated unique ways of exploration of dim-6 operators with specific focus on interference with the SM
- Gained access to unexplored VBS processes and observed them with full Run II dataset and improved data analysis methods
- The exploration of dim-8 operators carried out in myriad VBS and triboson topologies providing unique access to the full spate of possible new physics contributions
  - EFT validity constraints are taken into account in all cases
- Detailed exploration possible as we move towards Run III and beyond and improvements foreseen with machine learning based approaches (a la CMS ttZ EFT exploration)

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#### **Additional Material**

#### WW and WZ with two additional jets

#### CMS: Phys. Lett. B 809 (2020) 135710

#### Zγ + 2 jets



Access to neutral quartic vertex (like previous analysis)

Limits on T8 and T9 operators improve by a factor of two in comparison to ZZ + 2 jets process



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# Exploration of WW and WZ processes with two additional jets



Saptaparna Bhattacharya

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#### **Factoring in EFT validity**

- EFT validity taken into account by restricting EFT integration at the unitarity limit
- Unitarity limit is at ~1.5 TeV
- Calculated with VBF NLO (VBFNLO 1.4.0)

	Observed ( $W^{\pm}W^{\pm}$ )	Expected ( $W^{\pm}W^{\pm}$ )	Observed (WZ)	Expected (WZ)	Observed	Expected
	$(\text{TeV}^{-4})$	(TeV <sup>-4</sup> )	$(\text{TeV}^{-4})$	$(\text{TeV}^{-4})$	$(\text{TeV}^{-4})$	$(\text{TeV}^{-4})$
$f_{\rm T0}/\Lambda^4$	[-1.5, 2.3]	[-2.1, 2.7]	[-1.6, 1.9]	[-2.0, 2.2]	[-1.1, 1.6]	[-1.6, 2.0]
$f_{\rm T1}/\Lambda^4$	[-0.81, 1.2]	[-0.98, 1.4]	[-1.3, 1.5]	[-1.6, 1.8]	[-0.69, 0.97]	[-0.94, 1.3]
$f_{\rm T2}/\Lambda^4$	[-2.1, 4.4]	[-2.7, 5.3]	[-2.7, 3.4]	[-4.4, 5.5]	[-1.6, 3.1]	[-2.3, 3.8]
$f_{\rm M0}/\Lambda^4$	[-13, 16]	[-19, 18]	[-16, 16]	[-19, 19]	[-11, 12]	[-15, 15]
$f_{\rm M1}/\Lambda^4$	[-20, 19]	[-22, 25]	[-19, 20]	[-23, 24]	[-15, 14]	[-18, 20]
$f_{\rm M6}/\Lambda^4$	[-27, 32]	[-37, 37]	[-34, 33]	[-39, 39]	[-22, 25]	[-31, 30]
$f_{\rm M7}/\Lambda^4$	[-22, 24]	[-27, 25]	[-22, 22]	[-28, 28]	[-16, 18]	[-22, 21]
$f_{\rm S0}/\Lambda^4$	[-35, 36]	[-31, 31]	[-83, 85]	[-88, 91]	[-34, 35]	[-31, 31]
$f_{\rm S1}/\Lambda^4$	[-100, 120]	[-100, 110]	[-110, 110]	[-120, 130]	[-86, 99]	[-91, 97]

Limits worse by a factor of 5 in comparison with limits without unitarity constraints but more physical

#### ZZ with two additional jets

#### CMS: <u>Phys. Lett. B 812 (2020) 135992</u>

#### ZZ production with 2 jets



 Electroweak production of ZZ observed with 4.0 (3.5) σ observed (expected) significance

Provides access to the quartic coupling

Exploration of dim-8 operators performed

Provides access to T8 and T9 operators not possible in other final states

$$\mathcal{L}_{T,8} = B_{\mu\nu}B^{\mu\nu}B_{\alpha\beta}B^{\alpha\beta}$$
$$\mathcal{L}_{T,9} = B_{\alpha\mu}B^{\mu\beta}B_{\beta\nu}B^{\nu\alpha}$$

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#### ZZ production with 2 jets



Coupling	Exp. lower	Exp. upper	Obs. lower	Obs. upper	Unitarity bound
$f_{\rm T0}/\Lambda^4$	-0.37	0.35	-0.24	0.22	2.4
$f_{\rm T1}/\Lambda^4$	-0.49	0.49	-0.31	0.31	2.6
$f_{\rm T2}/\Lambda^4$	-0.98	0.95	-0.63	0.59	2.5
$f_{\rm T8}/\Lambda^4$	-0.68	0.68	-0.43	0.43	1.8
$f_{\rm T9}/\Lambda^4$	-1.5	1.5	-0.92	0.92	1.8

Access to neutral operators

Set strict bounds on these sets of operators, inaccessible in charged gauge boson final states

	Results	presented	factor in	unitarity	bound	S
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#### Wy with two additional jets

#### CMS: Phys. Lett. B 811 (2020) 135988

## Wy + 2 jets



First observation of the Wγ + 2 jets process with observed (expected) significance of 5.3 (4.8) σ

Exploration of dim-8 operators possible due to presence of SM quartic coupling

Invariant mass of the Wγ system is sensitive to presence of dim-8 operators

## Wy + 2 jets

Exploration of full set of "mixed" operators performed

 $\mathbf{M}$  For the parameters  $f_{M2-5}$ ,  $f_{M6-7}$  most stringent limits



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#### Zγ with two additional jets

#### CMS: <u>https://arxiv.org/abs/2106.11082</u>

#### Zγ + 2 jets



Access to neutral quartic vertex (like previous analysis)

Allows exploration of T8 and T9 operators

Invariant mass of the Zγ system is sensitive to presence of dim-8 operators

# Zγ + 2 jets



Coupling	Exp. lower	Exp. upper	Obs. lower	Obs. upper	Unitarity bound
$F_{\rm M0}/\Lambda^4$	-12.5	12.8	-15.8	16.0	1.3
$F_{\rm M1}/\Lambda^4$	-28.1	27.0	-35.0	34.7	1.5
$F_{\rm M2}/\Lambda^4$	-5.21	5.12	-6.55	6.49	1.5
$F_{\rm M3}/\Lambda^4$	-10.2	10.3	-13.0	13.0	1.8
$F_{\rm M4}/\Lambda^4$	-10.2	10.2	-13.0	12.7	1.7
$F_{\rm M5}/\Lambda^4$	-17.6	16.8	-22.2	21.3	1.7
$F_{\rm M7}/\Lambda^4$	-44.7	45.0	-56.6	55.9	1.6
$F_{\rm T0}/\Lambda^4$	-0.52	0.44	-0.64	0.57	1.9
$F_{\mathrm{T1}}/\Lambda^4$	-0.65	0.63	-0.81	0.90	2.0
$F_{\mathrm{T2}}/\Lambda^4$	-1.36	1.21	-1.68	1.54	1.9
$F_{\rm T5}/\Lambda^4$	-0.45	0.52	-0.58	0.64	2.2
$F_{\rm T6}/\Lambda^4$	-1.02	1.07	-1.30	1.33	2.0
$F_{\mathrm{T7}}/\Lambda^4$	-1.67	1.97	-2.15	2.43	2.2
$F_{\rm T8}/\Lambda^4$	-0.36	0.36	-0.47	0.47	1.8
$F_{T9}/\Lambda^4$	-0.72	0.72	-0.91	0.91	1.9

Stringent limits on dim-8 operators

Limits on T8 and T9 operators improve by a factor of two in comparison to ZZ + 2 jets process

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## Text book example of effective field theories



#### Special coordinate system



 $\mathbf{M}$  Define frame by a Lorentz boost to the center-of-mass frame of the W<sup>±</sup> $\gamma$  system

Solution Momenta are back-to-back

 $\mathbf{M}$  Angle  $\Phi$  now acquires sensitivity to interference with SM
## SMP-20-005: W±γ Results



 $\overline{\sigma(C_{3W})} = \sigma_{SM} + C_{3W}\sigma_{\text{interference}} + C_{3W}^2\sigma_{BSM}$ 

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### SMP-20-014: Measurement of the WZ process





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## **Power of new detectors**

- Machine learning with 3-D point clouds
- Extract information from edges of nearest neighbors
- Energy regression results are encouraging



# The Higgs Boson found at CERN!



# The CMS Experiment













# The Standard Model of Particle Physics

#### **Observations:**

- electron: 1897 (JJ Thomson)muon: 1936 (Anderson &
- Neddermeyer)
- electron neutrino: 1956
- (Cowan & Reines)
- •muon neutrino: 1962 (BNL)
- up, down, strange quark: 1968 (SLAC)
- charm quark: 1974 (SLAC/ BNL)
- tau lepton: 1975 (SLAC)
- bottom quark: 1977 (FNAL)
- gluon: 1979 (DESY)
- W and Z bosons: 1983 (CERN)
- top quark: 1995 (FNAL)
- tau neutrino: 2000 (FNAL)
- Higgs boson: 2012 (CERN)

## **Standard Model of Elementary Particles**



## **The Snowmass Process**

"The Particle Physics Community Planning Exercise ("Snowmass") is a process that takes place approximately every 6-8 years. Organized by the American Physical Society (APS) Division of Particles and Fields (DPF), Snowmass is an opportunity for the entire HEP community to plan and document a longterm vision for particle physics in the US along with its international partners

Thanks to the Community Planning Meeting committee

## WZ limits - interference and BSM separated



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# Salient features associated with VVV

 $\mathbf{Z}$ 

- Study of the Higgs sector possible
  - Best way to include Higgs mediated mode?
  - Can we quantify the interference between the Higgsmediated modes and other modes of production?
- Critical questions such as (dim6)<sup>2</sup> vs. dim8 contributions possible to address
  - Pertinent in the context of global EFT fits
- Explore sensitive variables: proxy for  $\widehat{S}$  and angular variables ( $\Delta \phi$ )
  - Absence of "golden" variables as in the case of vector boson scattering topologies



S

Operators	www	wwz	WZZ	ZZZ
$\mathcal{O}_{3\mathrm{W}} = \epsilon^{abc} W^{a\nu}_{\mu} W^{b\rho}_{\nu} W^{c\mu}_{\rho}$				
$\mathcal{O}_{\rm HD} = (D^{\mu}H^{\dagger}H)(H^{\dagger}D_{\mu}H)$				
$\mathcal{O}_{\rm HWB} = H^{\dagger} \sigma^{i} H W^{i}_{\mu\nu} B^{\mu\nu}$				
$\mathcal{O}_{\rm HW} = H^{\dagger} H W_{\mu\nu} W^{\mu\nu}$				

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# Relevant Operators for specific processes

Operators	wwww	WWZZ	ZZZZ	ZZZγ
$\mathscr{L}$ S,1 $\mathscr{L}$ S,2				0
ℒм,0 ℒм,1 ℒм,6 ℒм,7				
Шм,2 Шм,3 Шм,4 Дм,5	0			
$\mathscr{L}$ t,0 $\mathscr{L}$ t,1 $\mathscr{L}$ t,2				
$\mathscr{L}$ t,5 $\mathscr{L}$ t,6 $\mathscr{L}$ t,7	0			
${\mathscr L}$ т,8 ${\mathscr L}$ т,9	0	0		